

Experimental Study on Solid Circulation in a Multiple Jet Fluidized Bed

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Particle image velocimetry was used to investigate the evolution of multiple inlet gas jets located at the distributor base of a two-dimensional fluidized bed setup. Results were used to estimate the solid circulation rate of the fluidized bed as well as particle-entrainment into the individual jets. The effects of fluidization velocity, orifice diameter, orifice pitch, particle diameter, and particle density were studied. It was determined from this study that the solid circulation rate linearly increased with an increase in the fluidization velocity until the multiple jet system transitioned from isolated to an interacting system. In the interacting system of jets, the solid circulation increased with fluidization velocity but at a much lower rate. For multiple jet systems, this phenomenon may indicate the presence of an optimum operating condition with high circulation rate and low air input in the bed. © 2011 American Institute of Chemical Engineers AICHE J, 58: 3003–3015, 2012

Keywords: fluidization, multiple jets, particle image velocimetry, solid volume fraction, solid circulation, particle entrainment

Introduction

Fluidized beds commonly consist of multiple jets of oxygen containing gas located at the bottom of the bed to fluidize particulate media as well as support the combustion and gasification of solid fuel. The importance of these jets is reflected by their ability to act as a reacting medium and to control the reactors that produce syngas from gasification of fuels such as coal and biomass. These jets entrain gas and particles from the emulsion thereby enhancing the gas-particle mixing. Thus, the entrainment process helps to promote heat transfer, mass transfer, and chemical reaction rates for a fluidized bed reactor. For proper performance estimation, the patterns and extent of gas-particle mixing must be determined for a reactor. Failure to predict these factors may lead to problems such as particle agglomeration, temperature hot spots, and partial defluidization.

Various mathematical and empirical models for particle entrainment rate have been suggested in the literature. One of the earliest studies in this regard was done by Thorley et al.¹ using semicylindrical spouted beds. In the research conducted by Thorley et al.,¹ the forces acting on particles in the jet spout were balanced to estimate the particle velocity profile. Lefroy² suggested an alternative approach to predict the solid circulation rate by solving mass and momentum equations for bed voidage and particle velocities in the jet spout. This approach was further supplemented by various works in the literature.^{3–7} Van Deemter⁸ developed the counter-current flow model for gas–solid fluidized beds to propose the bubble induced particle mixing and solid circula-

tion. This was further extended by Baeyens et al.⁹ and van Deemter¹⁰ into comprehensive models and empirical equations to estimate the model parameters, particle velocities, circulation flux, bed turnover time, gas back mixing, residence time, segregation, and defluidization. Merry¹¹ developed a particle entrainment model based on particle stream function and validated it experimentally using a liquid–solid fluidized bed. Donsi et al.¹² developed a mathematical model based on the turbulent jet theory¹³ to describe the observed jet characteristics in a two-dimensional bed. This was further extended by Massimilla¹⁴ to a more detailed model which included mode of discharge, jetting regions, penetration depths, expansion angle, entrainment, bubble size and transport models. Further computational studies were conducted by various researchers to understand the influence of solid circulation on bed dynamics.^{5,10–14}

Various experimental studies have been performed to provide the empirical input parameters for these mathematical models, as well as to provide the data for model validation.^{3,15,16} The most important parameters to quantify the particle entrainment rate are particle velocities and void fraction values. These parameters are estimated either inside the jet or along the jet periphery. Several intrusive and non intrusive techniques have been applied to quantify these parameters in experimental fluidized beds setups. Some of these techniques include photographs and visual analysis,^{1,11,12,17–20} pitot tubes,^{12,18,21} fiber optic probes,^{4,22–27} tracer particles,^{7,21,28,29} tracer gas,^{19,30} gamma-ray densitometer,^{15,31} positron emission particle tracking,^{32,33} x-ray tomography,^{34,35} digital image analysis,^{36–39} and particle-image velocimetry (PIV).^{40,41} Most of these experimental studies have been conducted using single jet spouted beds. Also, reports on the influence of multiple jet systems on solid circulation are still sparse in the literature.

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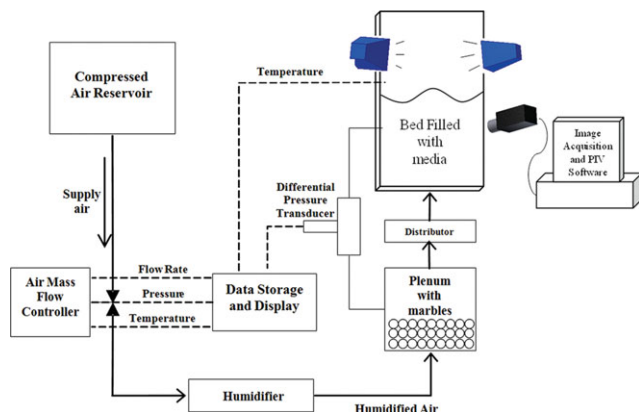


Figure 1. Schematic of experimental setup.

[Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

The motivation behind this study was to understand the effect of distributor orifice diameter, orifice pitch, particle types, and operating conditions on the solid circulation dynamics of a multiple jet system fluidized bed. PIV has recently been used on a single jet system to analyze the solid circulation behavior.⁴⁰ However, its application onto multiple jet systems has not been reported in the literature. In this study, PIV and a two-dimensional bed were used to visualize and analyze the solid circulation phenomena, including jet entrainment, on a multiple jet system.

Experimental

Experimental setup

A two-dimensional fluidized bed setup of dimensions 0.76 m high and 0.25 m wide was used in this research. Figure 1 contains a schematic of the experimental setup. The thickness of the bed was 0.0127 m to ensure that the ratio of bed thickness to particle diameter is greater than 10 to avoid particle bridging effects.^{42,43} The bed walls were made of high optical clarity glass, and a light absorbing black sheet was located on the back side. The bed was illuminated homogeneously using two halogen lamps of 500 W each, which were located approximately at a 45° angle from the bed axis. A Flowsense MKII (4 Megapixel) PIV camera was used to take double frame digital images of the bed. The camera was controlled through a timer box, and images were stored and processed in a computer using Dynamic Studio 3.0 software. The setup was surrounded by black barriers and curtains to avoid reflections in captured images.

Air supply was controlled through a 500 SLPM ALICAT (Model No. MCR-500SLPM-D/10M) mass flow controller (accuracy: $\pm 0.4\%$ reading, $+0.2\%$ FS). As shown in Figure 1, the air was channeled through a humidifier to avoid electrostatic effects in the bed. The water content of the humidifier was kept to a minimum to avoid saturation. Air was then fed into the plenum chamber and injected into the bed through a perforated plate distributor. The plenum was filled

Table 1. Physical Properties of Tested Fluidized Bed Media

Particle	d_p (μm)	ρ_p (kg/m^3)	U_{mf} (m/s)	Geldart Group
Glass	550	2500	0.22	B
Glass	750	2500	0.35	D
Ceramic	550	1079	0.20	B

Table 2. Description of Perforated Distributor Plates

Distributor Type	d_o (mm)	N (Number of Holes)	P (m)
A1	2	9	0.028
A2	2	5	0.056
B1	3	4	0.063

with marbles to ensure uniform fluidization. A Setra (Model no. 264) differential pressure transducer (ranging 0–25,000 Pa, accuracy $\pm 0.25\%$ FS) was used to measure the pressure drop between the plenum and the freeboard. Differential pressure, inlet pressure, inlet temperature and air flow rate data were recorded using a National Instruments data acquisition system.

Test cases

Test cases were developed to explore the effect of changing particle size, particle density, fluidization velocity, orifice diameter and orifice pitch on the solid circulation behavior. Three types of particles were selected to provide diameter ranging 550–750 μm and density ranging 1079–2500 kg/m^3 . As listed in Table 1, this resulted in particles with Geldart B and D classification. Ceramic 550 μm and glass 550 μm particles have the same mean diameter, different density, and the same Geldart B classification. Glass 550 μm and glass 750 μm particles differ in mean diameter and Geldart classification but have same density. The particles were sieved through calibrated sieves multiple times within 100 μm range. The particle-size distribution was assumed to be normal and the mean diameter is reported in Table 1. Each type of particle was examined using three different types of perforated distributor plates. As listed in Table 2, effects of orifice diameter and orifice pitch were evaluated in this study. Distributor Type A1 and A2 have same orifice diameter but different orifice pitch. Distributor Types A and B differ in orifice diameter. Distributor Type A1 and B1 were designed to have the same total orifice flow area. Thus, at same superficial velocity for Type A1 and B1 distributors, the effect of orifice diameter can be compared at the same orifice velocity. The uncertainty in orifice diameter was within ± 0.02 mm to ensure the flow rate uniformity. The periphery of the orifices is shown in Figure 2 to be located 1.59 mm from the bed wall aiding in the optical assessment of jets. Similar orifice arrangements have been used for optical assessment

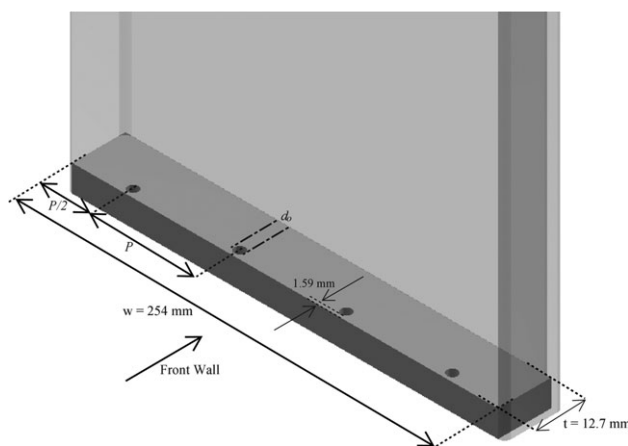


Figure 2. Schematic of distributor plate.

Table 3. Comparison of Experimental Minimum Fluidization Velocities Measured for Cylindrical and Two-Dimensional Bed Test Runs

Distributor Type	Glass 550 μm Particles (m/s)	Glass 750 μm Particles (m/s)	Ceramic 550 μm Particles (m/s)
3D Bed ^a	0.22	0.35	0.20
A1 ^b	0.23 [1.0]	0.34 [1.0]	0.23 [1.1]
A2 ^b	0.21 [0.9]	0.28 [0.8]	0.17 [0.9]
B1 ^b	0.23 [1.0]	0.34 [1.0]	0.18 [0.9]

Term in bracket indicates $U_{\text{mf},2\text{D}}/U_{\text{mf},3\text{D}}$ for the corresponding particle type.

^a $U_{\text{mf},3\text{D}}$ (m/s).

^b $U_{\text{mf},2\text{D}}$ (m/s).

of jet characteristics in two-dimensional and semicylindrical beds with a flat front wall.^{12,17,21,44–50}

Recent review articles have cautioned the application of two-dimensional fluidized bed data for quantitatively predicting cylindrical fluidized beds.^{42,43} However, it was acknowledged that the two-dimensional beds provide qualitative and semiquantitative indications of behavior for an axis-symmetric cylindrical column. It has also been recommended^{42,43} to keep the ratio of bed thickness to particle diameter high (approximately 10–20) to avoid any particle bridging effects in a two-dimensional fluidized bed setup. In absence of such conditions, frictional resistance becomes a dominant force in restraining the particle motion along the front and rear faces of the bed.⁴³ In the current study, the ratio of bed thickness to particle diameter was kept at 23 for glass and ceramic 550 μm particles and 17 for glass 750 μm particles.

Analysis Procedure

Pressure drop analysis

The differential pressure drop was analyzed to calculate the minimum fluidization velocity ($U_{\text{mf},2\text{D}}$) for each combination of particle type and distributor type in the two-dimensional bed. These velocities are listed in Table 3 and compared with minimum fluidization velocities ($U_{\text{mf},3\text{D}}$) calculated for respective particles in a cylindrical fluidized bed setup.⁵¹ Experimental data was taken starting with the bed at the maximum flow rate, which was approximately two times the minimum fluidization velocity. The flow rate was decreased by 5 SLPM and then held constant for 300 s before the data acquisition.

Particle image velocimetry

PIV is a nonintrusive optical technique used to quantify velocity field information of seeding particles suspended in a fluid through measurement of particle displacement over a known time interval. Due to dense media prohibiting laser passage in the bed, a conventional PIV technique cannot be used to obtain gas or particle velocities. However, PIV can be used to quantify particle velocities near the front wall. Hence, the bed was homogeneously illuminated using halogen lights and bed media was used as seeding particles to quantify the particle velocity field.

The camera was located at a distance of 1.0 m from the front wall which enabled a field-of-view covering the full width and expanded height of the bed. Magnification factor of this setup ensures that each particle occupies approximately 3–4 pixels on a CCD frame of the 4 Megapixel PIV camera. Measurements at each flow rate included 25 double frame images taken at a trigger rate of 7.4 Hz. This resulted in measuring the flow field over 3.38 s. The time delay

between each pair of double frame images was 1 ms. An adaptive correlation technique was applied where the interrogation area (IA) for each reading was iteratively reduced from 128×128 pixels to 32×32 pixels in two steps. The relative overlap of 25% and the peak validation of 1.2 were applied for PIV analysis. A sample PIV mean superimposed onto the mean image is shown in Figure 3. PIV vectors inside the jet were masked, as the time delay was set to get better velocity profiles in the annulus area outside the jets.

Solid volume fraction analysis procedure

Quantifying the solid volume fraction, ϵ_s , of a fluidized bed is necessary to investigate its hydrodynamics. In two-dimensional fluidized beds, the images of fluidizing particles can be obtained at the bed front wall. The time average intensity analysis of these images can be used to determine the solid volume fraction in a two-dimensional bed.^{36,38,39,52–54} This section discusses the procedure used to determine the solid volume fraction values for the bed.

Corrected Intensity Maps. Twelve bit images were captured in these experiments, ranging the pixel intensity level from 0 (for dark regions) to 4095 (for bright regions) for a window of 2048×2048 pixels. Fluidizing bed images were subtracted by a base image to obtain a homogeneously lit image. The base image was taken without any particles in the bed to remove the illumination irregularities. The resultant image intensities were then scaled from 0 (absolute gas phase) to 1 (particulate phase). This scaling ensured that different type particles (glass or ceramic) with varying spectral properties are equally treated mathematically. Thereafter, the scaled intensity of each pixel was spatially averaged for the same size as the PIV vector Interrogation Area (IA). This enabled each of the 85×85 PIV velocity vectors in the image window to have a corresponding averaged intensity value. Lastly, the images taken at different air flow rates were linearly corrected using the intensity difference caused by the level of smearing in those images due to moving particles to get the corrected intensity, I_{cor} , for each IA. This image analysis is shown for an instantaneous image in Figure 4. The image processing work was performed using MATLAB 7.9 (R2009b).

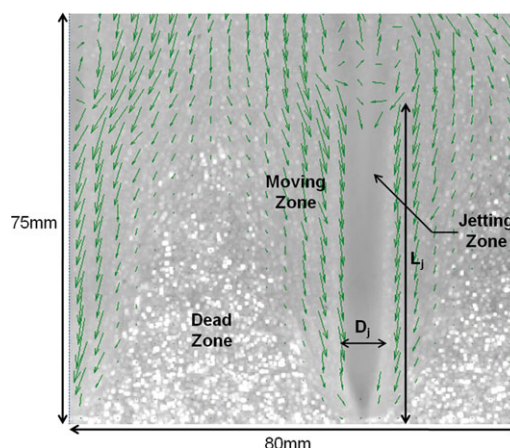


Figure 3. Jet at the bottom of Type A2 distributor plate using glass 550 μm particles: PIV mean superimposed onto mean image of acquired images.

[Color figure can be viewed in the online issue, which is available at [wileyonlinelibrary.com](http://www.wileyonlinelibrary.com).]

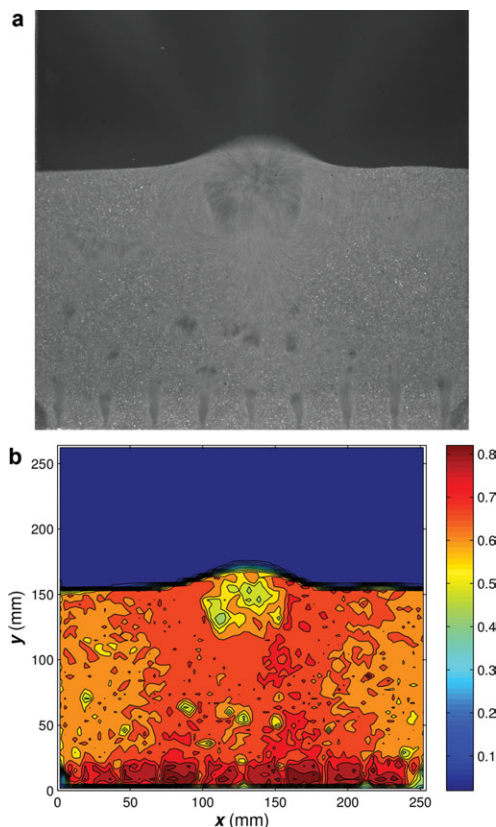


Figure 4. Image analysis for the test case of glass 550 μm particles using Type A1 distributor at $U_f = 0.31$ m/s: (a) Instantaneous image with non-homogenous illumination (b) Corrected intensity, I_{con} map of the instantaneous image.

0 represents all gas and 1 represents all particulate phase. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

Bulk Solid Fraction. The bulk solid volume fraction, $\varepsilon_{s,b}$, of the fluidized bed was calculated at a set fluidization velocity to develop a solid volume fraction correlation based on I_{cor} of the respective IAs. For this purpose, the corrected intensity, I_{cor} , maps were time-averaged for the data acquisition period of 3.38 s. These time-averaged I_{cor} maps were used to quantify the bulk solid volume fraction, $\varepsilon_{s,b}$, of the bed at a set air flow rate. First, the time-averaged I_{cor} intensity value was averaged along the bed width and plotted along the height above distributor, as shown in Figure 5a. The bed expansion height was the vertical location at which this intensity is 5% of the maximum value (1.0). The calculated bed expansion height was subsequently used to determine the $\varepsilon_{s,b}$ using the known bed mass, as shown in Figure 5b.

Solid Volume Fraction. The time-averaged I_{cor} value for each interrogation area was converted to the mean solid volume fraction, ε_s , along the thickness of the bed using the following correlation.

$$\varepsilon_s = (AI_{\text{cor}})^B \quad (1)$$

The term AI_{cor} represents the area fraction occupied by illuminated particles on the front wall of the fluidized bed. This includes higher intensity particles at the front wall and

lower intensity particles along the thickness of the bed. The basis for the assumption that the measured intensity represents the particle area fraction came from the studies conducted by Heffels et al.⁵⁵ and Boerefijn et al.,⁵³ in which they concluded that the backscattered light reflected from the particle suspensions is an indication of the total reflecting surface area of the particles. A similar conclusion was made by Poletto et al.⁵² using image analysis on a water fluidized bed of $3d_p$ bed thickness.

The parameter B was introduced in Eq. 1 to convert this area fraction to mean solid volume fraction, ε_s , along the thickness of the bed. This was followed by the work of Link et al.,^{36,37} where porous cube theory was applied onto two-dimensional fluidized beds to obtain a correlation between its intensity and solid volume fraction

$$\varepsilon_s = (I_m)^{1.5} \quad (2)$$

In the correlation proposed by Link et al. in Eq. 2, I_m represents the area fraction of particles on the front wall computed through intensity manipulation, which is similar to AI_{cor} term in Eq. 1 with $A = 1$. I_m was raised by a power of 1.5 to obtain mean ε_s along their bed thickness ($6d_p$). This

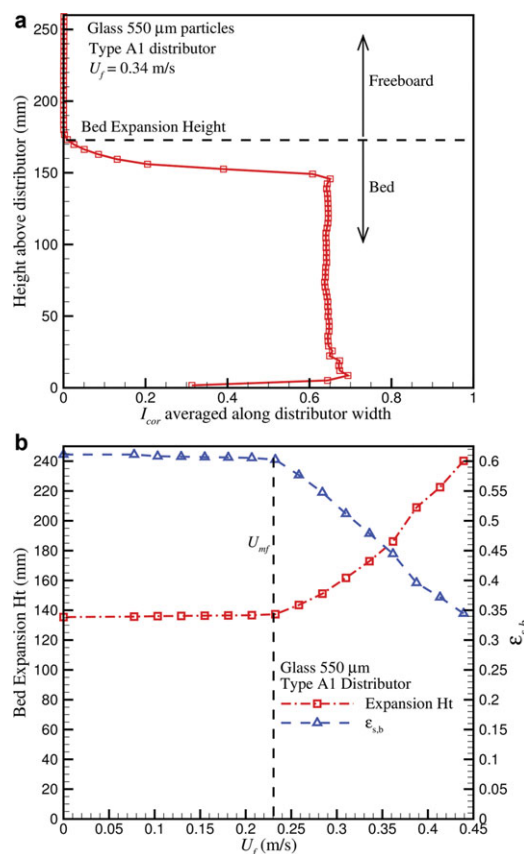


Figure 5. Bulk solid volume fraction, $\varepsilon_{s,b}$, estimation for glass 550 μm particles using Type A1 distributor.

Bed mass is 0.6674 kg: (a) Corrected Intensity averaged along distributor width vs. height above distributor at $U_f = 0.35$ m/s (b) Bed expansion height and $\varepsilon_{s,b}$ for bed superficial velocities. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

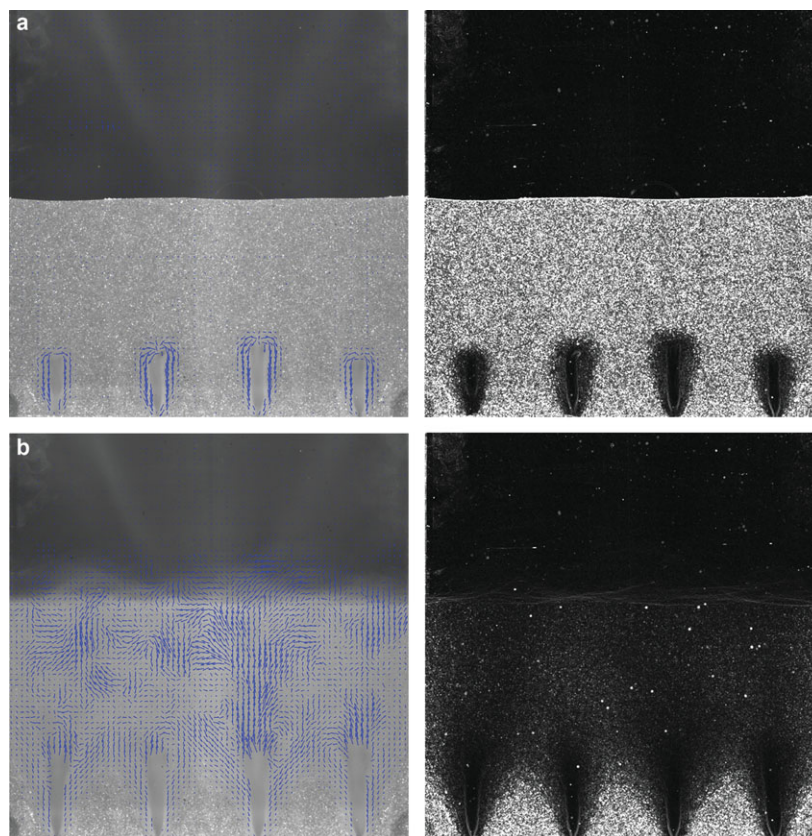


Figure 7. PIV data superimposed on mean image of acquired images, and Prewitt filter applied on mean image.

Test-case of Type B1 distributor using glass 550 μm particles exhibiting different jet systems (a) $U_f/U_{mf,2D} = 1.12$, isolated system and (b) $U_f/U_{mf,2D} = 1.46$, interacting system. Field of view is 250 mm wide by 200 mm high. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

Figure 8b provides the plot for absolute error in total bed mass calculated using the solid fraction model, Eq. 1, against the actual mass of particles fluidized in the bed. This error was found to be less than 15% for all nine test cases at varying fluidization velocities. Figure 8c contains the void fraction map ($\epsilon_f = 1 - \epsilon_s$), calculated using Eq. 1, for the instantaneous image shown in Figure 4a.

Figure 9a contains the plot of solid volume fraction profile along the jet boundary (a-b-c), including the conical bottom, for ceramic 550 μm particles using Type A1 distributor plate at $U_f/U_{mf,2D} = 1.46$. The solid fraction values increase almost linearly along the cone length to attain its maximum value at the location where conical part of the jet structure ends (a-b). This indicates the presence of dead zones in the vicinity of conical regions. Lower solid fraction near the bottom part of the jet is attributed to higher diffusion rate of high velocity gas from the jet into the bed. Further upwards along the jet height (b-c), the solid fraction profile gradually decreases along the constant diameter part of the jet indicating the influence of the particle moving zone on the solid fraction values. Figure 9a also contains the solid volume fraction profile along the center of jet 2 (p-q). The solid fraction values at the center are lower in comparison to the values at the boundary of the corresponding jet. The solid fraction values at the center follow the profile similar to that at the boundary, except for the upper part of the jet where a gradual increase in solid fraction was measured. This complements the observation at the jet boundary where the entrainment of particles leading to the decrease of solid fraction values now contributes to the increase of solid fraction

values at the center of jet. Similar results were observed for other test cases involving different distributor types and particle combinations.

The profiles shown in Figure 9a match qualitatively with the theoretical results of Donsi et al.¹² for solid fraction values along the jet central axis, and computational results of Ettehadieh et al.⁵ for solid fraction values 2.5 mm off the jet central axis. Quantitatively, the solid fraction values measured by Donsi et al.¹² and Ettehadieh et al.⁵ were lower than those observed in Figure 9a. The reason for this difference could partially be attributed to mean jet orifice velocity, U_o , in the aforementioned studies being ~ 2 times higher than the current study. Another comprehensive experimental study conducted by the same researchers,¹⁵ using rectangular slit distributors and comparable inlet jet velocities revealed the solid volume fraction values to be in the same range as observed within the current study.

Figure 9b contains solid volume fraction profiles at different horizontal positions along the height of jet. As expected, the solid fraction values at the center of the jet are lower than those at the jet boundary and further lower than those located within the dead zone ($x/P = \pm 0.5$). The solid fraction profile at the jet center follows the same pattern as shown for the second jet in Figure 9a. The solid fraction value at the center is highest at $y = 0.010$ m and it gradually increases along the jet height from $y = 0.020$ m to $y = 0.060$ m. At the height of $y = 0.060$ m, the solid fraction values are almost uniform along the pitch of the distributor indicating the end of the jet length. The solid fraction values at the two jet boundaries are averaged and represented in Figure 9a. This

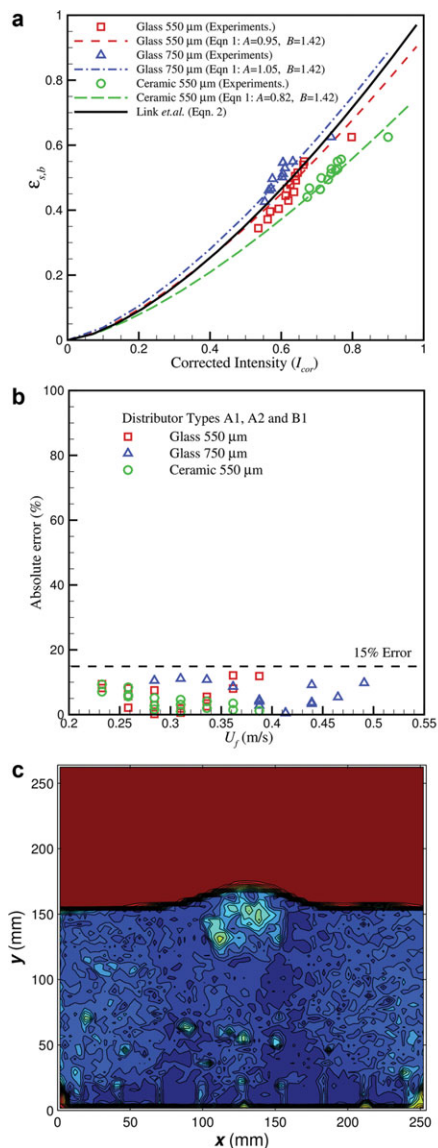


Figure 8. Validation and application of solid volume fraction correlation.

(a) Comparison of ϵ_s correlation (Eq. 1) with correlation from Eq. 2 and $\epsilon_{s,b}$ values for different test cases. (b) Absolute error between actual bed mass and the mass calculated using ϵ_s correlation (Eq. 1) along the expanded bed dimensions (c) Void fraction ($\epsilon_f = 1 - \epsilon_s$) map calculated using Eq. 1 for the instantaneous image shown in Figure 4. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

average value decreases slightly for height greater than 0.020 m due to the entrainment of particles from the moving zone into the jetting zone. The solid fraction values at the center of dead zones ($x/P = \pm 0.5$) do not show any change along the jet length for heights greater than 0.020 m. This horizontal profile of solid volume fraction was found to be similar to those seen in the literature.⁶¹

Particle entrainment and solid circulation

This section reviews the effect of variations in fluidization velocity, particle properties and distributor types on solid circulation and particle entrainment rates.

Fluidization Velocity. Figure 10a provides a plot of solid circulation rate for three particle types using the Type A1

distributor. Solid circulation was found to increase linearly with an increase in fluidization velocity until the jet systems transitioned from isolated to interacting jets. For the interacting jet systems, the rate of increase of solid circulation suddenly decreased in comparison to isolated jet systems. Although the linear increase of solid circulation with fluidization velocity has been observed in previous studies,^{20,21,28,29} the sudden change in the rate of solid circulation has not been reported in the literature due to the lack of multiple jet studies. The relative leveling of solid circulation rate with fluidization velocity indicates that an optimum operating condition may exist for the applications that require high solid circulation for lower air input flow rate.

The accuracy of these solid circulation rate values were evaluated by comparing them with an alternate method presented by van Deemter.⁸ The circulation rate was proposed to be calculated using a counter-current model developed for gas–solid fluidized beds^{62,6,3}

$$\dot{m}_{x,tot} = \rho_p \times (\epsilon_{s,b}) \times (fr_{neg}) \times (|v_{neg,avg}|) \times (A_d) \quad (3)$$

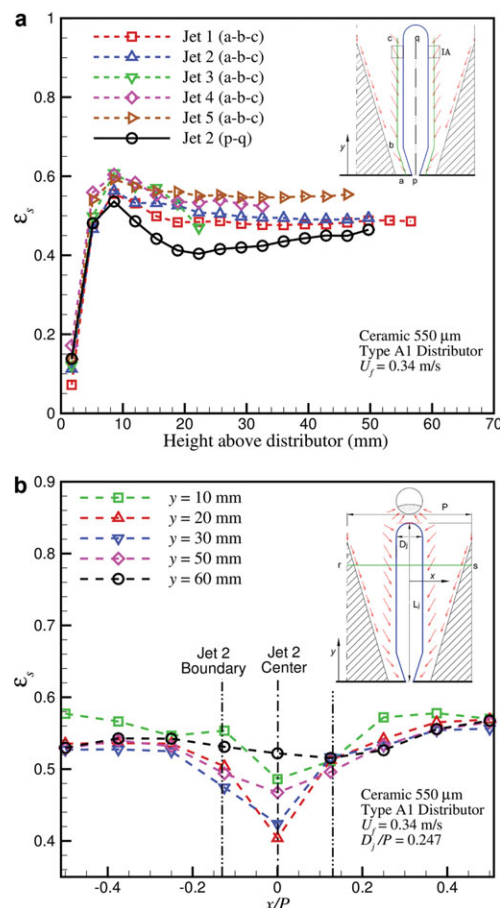


Figure 9. Solid volume fraction profiles for ceramic 550 μm particles with Type A1 distributor and $U_f = 0.34$ m/s.

(a) Solid volume fraction profile measured along jet boundary (a-b-c) and jet center (p-q) vs. height above distributor. Plot shows jet boundary data for five jets counting from left to right and jet center data for second orifice from the left. (b) Solid volume fraction profile for different horizontal locations (r-s) centered on the second jet from the left. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

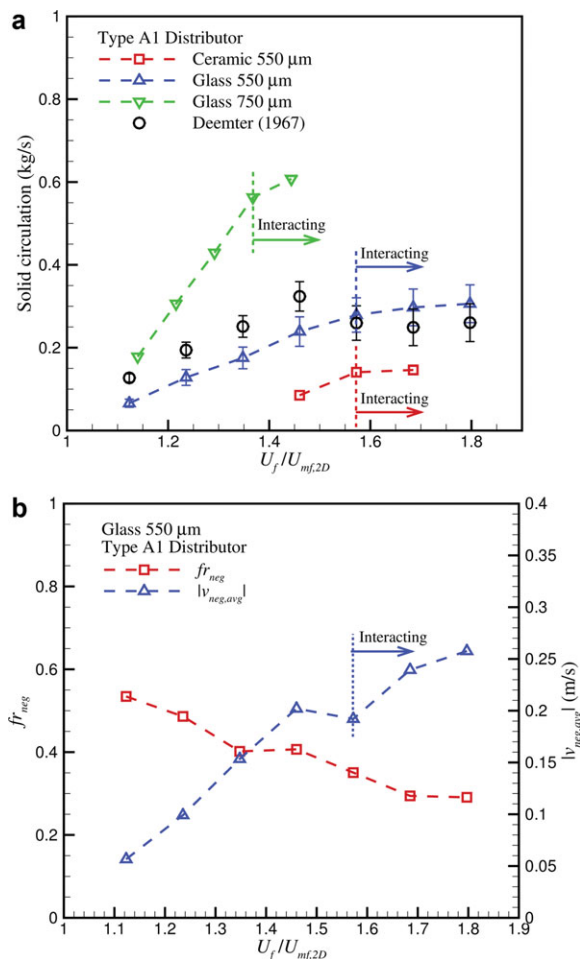


Figure 10. Analysis of solid circulation rate with change in fluidization velocity.

(a) Solid circulation rate calculated for the three particle types using Type A1 distributor vs. $U_f/U_{mf,2D}$. (b) $f_{r,neg}$ and $|v_{neg,avg}|$ calculated for glass 550 μm particles using Type A1 distributor vs. $U_f/U_{mf,2D}$. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

In Eq. 3, $\varepsilon_{s,b}$ represents the bulk solid volume fraction calculated by Eq. 1, $f_{r,neg}$ represents the fraction of dense bed moving downwards, $|v_{neg,avg}|$ represents the average negative particle velocity of the bed, and A_d represents the cross-sectional distributor area. These parameters were calculated for the test-case of glass 550 μm particles using the Type A1 distributor plate at different fluidization velocities. The calculated solid circulation rates are plotted in Figure 10a along with the average uncertainty estimation of $\pm 18.2\%$. These circulation rate values are close in comparison to the values calculated using the method suggested in the current study, for which the average uncertainty was estimated to be $\pm 14.9\%$. As observed from Figure 10a, the values calculated by Eq. 3 overestimated the current results for isolated/transitional jet systems and under predicted values for the interacting jet systems.

It must be noted that the PIV camera settings were adjusted to get better signal/noise ratio for particle velocities participating in the jet entrainment. Particle velocities near the top of the bed are higher due to the eruption of bubbles.^{64,65} This might lead to under prediction of $|v_{neg,avg}|$ using the current experimental settings, especially at higher

fluidization velocity. Also, there is a level of unaccounted error within the estimation of $f_{r,neg}$ which is more accurately noted as the fraction of downward moving particle and interstitial gases.¹⁰ It was not possible to accurately implement this definition onto the current results of the experimental study. However, the comparison between the data presented in Figure 10a shows reasonable agreement in solid circulation rate values calculated using the two methods. Figure 10b shows the parameters $f_{r,neg}$ and $|v_{neg,avg}|$ estimated for the discussed test-case. As noted, the $f_{r,neg}$ values were found to monotonically decrease with the increase in fluidization velocity. This trend can be partially explained on the basis of reduction of the dead zone areas and more particles participating in the fluidization of the bed due to the high velocity upward jets. The $|v_{neg,avg}|$ values were found to linearly increase with an increase in fluidization velocity, until the jet systems transitioned to interacting jets. For the interacting jet systems, a sudden decrease in $|v_{neg,avg}|$ values was observed which is consistent with the leveling of solid circulation rate observed in Figure 10a.

It was found that the major factor contributing to the leveling of solid circulation rate, calculated by the current experimental scheme, was the leveling of the horizontal particle entrainment velocity, \bar{u} , for the interacting jet systems. Figure 11 contains a plot of average entrainment velocities, U_{avg} , for the three test cases mentioned in Figure 10. U_{avg} was calculated while measuring the horizontal entrainment velocities along the jet boundaries and then averaging these values for all jets at a particular air flow rate, as shown in Figure 6. As observed from the Figure 11, U_{avg} increased linearly with the fluidization velocity until the jet systems transitioned to the interacting system, where U_{avg} started to level off. Since the jet length and jet diameter increased with fluidization velocity,⁵⁶ the leveling of horizontal entrainment velocities was the determining factor in the leveling off of the solid circulation rate.

Figure 12 provides the particle entrainment rate for each of nine individual jets with glass 550 μm particles and Type A1 distributor plate. It was observed that the particle entrainment in each jet is not same at a particular fluidization velocity. This difference in particle entrainment is primarily

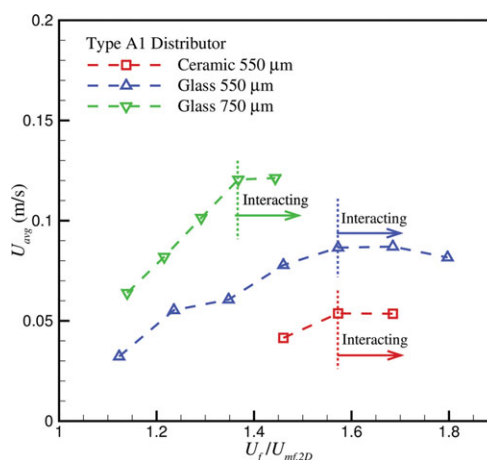


Figure 11. Average entrainment velocity, U_{avg} , vs. $U_f/U_{mf,2D}$ for three particle types using Type A1 distributor.

[Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

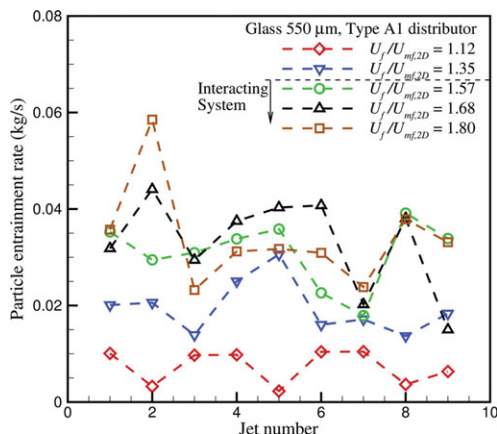


Figure 12. Particle entrainment rate for individual jets with glass 550 μm particles and Type A1 distributor plate.

Plot shows data for all 9 jets counting from left to right along the distributor width. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

attributed to differences in jet lengths and entrainment velocities for each jet. Jet diameter and solid volume fraction values along the jet periphery did not differ considerably for each jet, as shown in Figure 9a.

Data in Figure 12 demonstrate that the particle entrainment for a particular orifice does not necessarily increase with an increase in fluidization velocity. This phenomena was particularly evident for the interacting jet systems ($U_f/U_{mf,2D} = 1.57 - 1.80$). However, as noted in Figure 10a, a net increase in solid circulation $\dot{m}_{x,tot} = \sum_{j=1}^N \dot{m}_{x,i}$ is always observed with an increase in superficial velocity.

It was also noted that no particular orifice consistently had a maximum or minimum particle entrainment rate for all fluidization velocities. Hence, the particle entrainment rates were not biased for any particular orifice location. Aforementioned observations stress the differences between the current study and single jet experiments conducted in the literature, where such phenomena could not be noticed.

Particle Properties. Figure 13 contains a comparison of volumetric circulation rate for glass 550 μm and ceramic 550 μm particles using Type A1 distributor plate. The volumetric circulation rate was calculated by dividing the respective solid circulation rate of particle types with the particle density. As observed from the plot, the circulation rate of ceramic 550 μm particles was 14–17% higher than glass 550 μm particles at $U_f/U_{mf,2D} > 1.57$. For a relatively high difference in density of the mentioned particles, the difference in volumetric circulation rate is small. As observed from Figure 11, the average horizontal particle velocity, U_{avg} , along the jet periphery was higher for glass 550 μm particles. However, the ceramic 550 μm particle test cases produced higher jet lengths than glass 550 μm test cases, as shown in Figure 14. The jet diameter was not dependent on particle properties,⁵⁶ and hence did not significantly impact the differences in the circulation rates for the aforementioned particle types. The difference in the mean solid volume fraction, ϵ_s , for the two particle types was not found to be significant, with the profile for glass 550 μm particles being similar to that shown in Figure 9 for ceramic 550 μm particles. The relatively low number of data points for comparison of these two particle types is due to the fluidizing nature of ce-

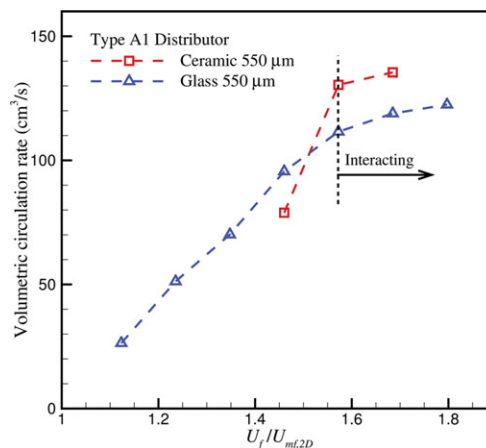


Figure 13. Volumetric circulation rate vs. $U_f/U_{mf,2D}$ for glass 550 μm and ceramic 550 μm particles using Type A1 distributor plate.

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ramic 550 μm particles. For higher fluidization velocities, jets for ceramic 550 μm particles bypassed the bed height and thus their jet length was not defined. At low fluidization velocities, the estimation of accurate jet diameter for ceramic 550 μm particles was not possible due to the instability of jets.

Figure 15 contains the comparison for solid circulation rate of glass 550 μm and glass 750 μm particles using Type B1 distributor plate. For glass 750 μm particles, the data was obtained for velocities both below and above its minimum fluidization velocity ($U_f/U_{mf,2D} = 1$). For glass 550 μm particles, the solid circulation rate was obtained only for velocities greater than its minimum fluidization velocity. As observed from the plot, the solid circulation rate for glass 750 μm particles exceeded the solid circulation rate of the glass 550 μm particles at $U_f/U_{mf,2D}$ greater than 1.

A similar trend was seen when both particles were tested with the Type A1 distributor plate, as shown in Figure 10a. This plot contains the data at velocity higher than minimum fluidization ($U_f/U_{mf,2D} = 1$) for both particle types. In this

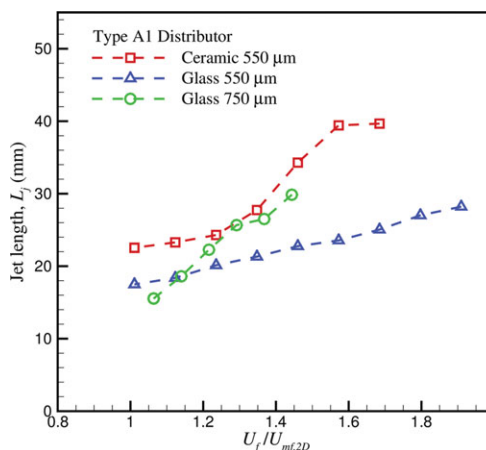


Figure 14. Jet penetration length vs. $U_f/U_{mf,2D}$ for the three particle types using Type A1 distributor plate.

[Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

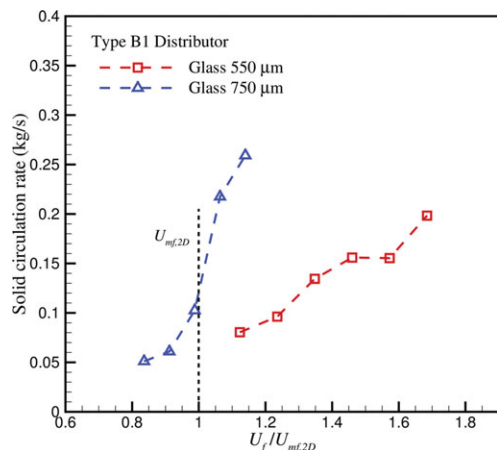


Figure 15. Solid circulation rate vs. $U_f/U_{mf,2D}$ for glass 550 μm and glass 750 μm particles using Type B1 distributor plate.

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comparison, again at $U_f/U_{mf,2D}$ greater than 1, the solid circulation rate for glass 750 μm particles was found to be higher than that of glass 550 μm particles. The jet penetration length for glass 750 μm particles was found to be higher than that of glass 550 μm particles, as seen from Figure 14. Also, as observed from Figure 11, the average velocity along the jet periphery, U_{avg} , was found to be higher for glass 750 μm particles.

In their respective single jet experiments, Filla et al.¹⁸ and Kim et al.²⁰ noticed an increase in entrainment rate for the increase in particle diameter. Filla et al.¹⁸ attributed both higher solid fraction values and higher entrainment velocities throughout the jet, leading to higher momentum exchange between the solids and gas, for the increase in entrainment rate for coarser particles. However, the theoretical model of Baeyens et al.⁹ disagreed with this impact of particle diameter on the particle entrainment values, where it was concluded that at equal values of excess velocity ($U_f - U_{mf}$) the entrainment rate decreases for coarser particles. José et al.²⁷ observed an increase in solid circulation rate for the decrease in particle density, and also acknowledged that the horizontal component of particle entrainment velocity decreases with the decrease in particle density. Filla et al.¹⁸ concluded that the resistance of the particulate phase to the dilute region of jet is larger for smaller and lighter particles, leading to lower entrainment velocities. These data agree with the findings presented in the current study.

Distributor Configuration. Distributor Type A1 and A2 were tested with glass 550 μm particles to understand the effect of the change in orifice pitch on the solid circulation rate. Figures 16a, b shows the solid circulation rate for aforementioned distributors vs.

$U_f/U_{mf,2D}$ and orifice velocity, respectively. As observed from Figure 16a, the solid circulation rate of Type A2 distributor was slightly higher, if not equal, in comparison to the data for Type A1 distributor at the same $U_f/U_{mf,2D}$. Due to fewer holes in Type A2 distributor, the orifice velocity for Type A2 is calculated to be higher than Type A1 at same fluidization velocities. Hence, as observed from Figure 16b, when the solid circulation data was observed for the same orifice velocities, the solid circulation rate for Type A1 distributor was found to be higher than Type A2. The solid cir-

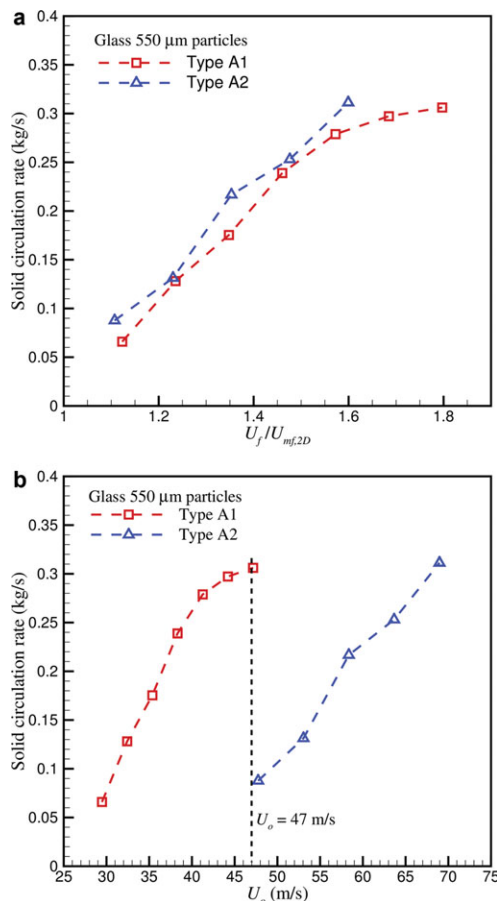


Figure 16. Comparison of solid circulation rates for Type A1 and A2 distributor plates using glass 550 μm particles.

(a) vs. $U_f/U_{mf,2D}$. (b) vs. orifice velocity. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

ulation for Type A1 distributor is much higher than Type A2 distributor due to higher average entrainment velocities, U_{avg} , and higher jet lengths⁵⁶ at the same orifice velocities. Although the particle entrainment rate measured for each jet

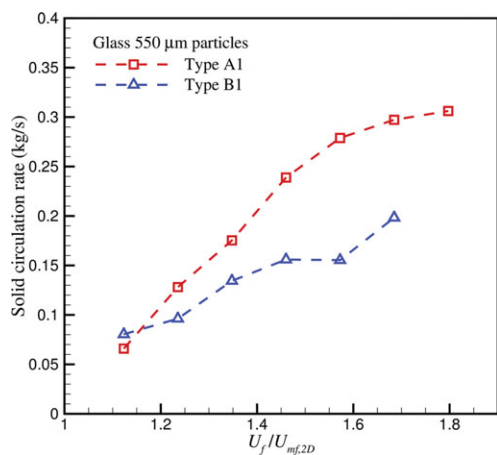


Figure 17. Solid circulation rate vs. $U_f/U_{mf,2D}$ for Type A1 and B1 distributor plates using glass 550 μm particles.

[Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

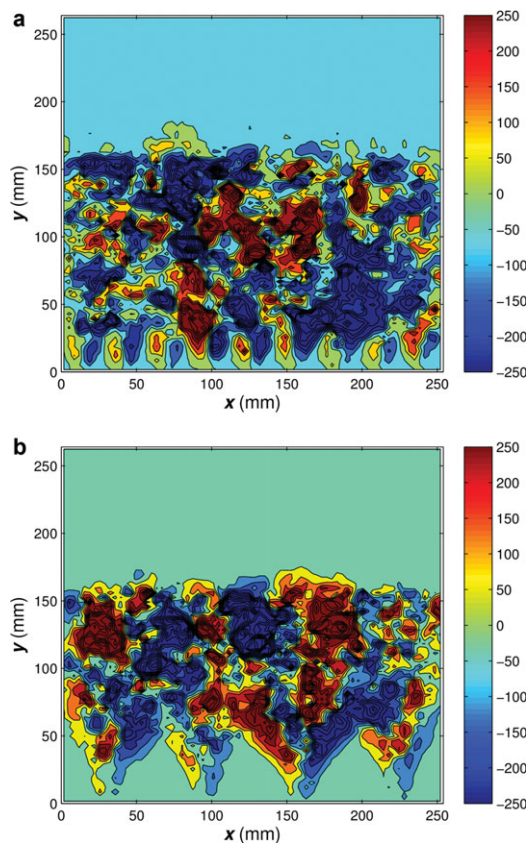


Figure 18. Comparison of horizontal particle mass flux maps, \dot{m}_x'' , at $U_f/U_{mf,2D} = 1.46$ using glass 550 μm particles (a) Type A1 distributor (b) Type B1 distributor.

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in Type A2 distributor plate was higher, the number of orifices in Type A2 was lower in comparison to Type A1 distributor plate. Therefore, the total solid circulation rate for the bed decreased for a distributor plate with higher orifice pitch when compared at the same orifice velocity.

The effect of the increase in orifice diameter on solid circulation rate was measured by testing glass 550 μm particles with distributor Types A1 and B1. Distributor Type A1 and B1 were designed to have the same total orifice flow area. Hence, at the same superficial velocity, the effect of orifice diameter can be compared at the same orifice velocity for these distributors. For Type B1 distributor, the jet penetration length⁵⁶ and particle entrainment rate per jet was found to be higher than Type A1 distributor. However, the average particle entrainment velocities and number of orifices were higher for Type A1 distributor. Figure 17 contains the solid circulation rate data for the two test cases. It should be noted from Table 3 that $U_{mf,2D}$ for both test-cases is same. As observed, the circulation rate for Type A1 distributor was measured to be 30–80% higher than Type B1 distributor plate. Hence, for a distributor plate design, the plate with large number of small diameter orifices was found to be effective at enhancing the circulation rate. This is also evident from the horizontal particle mass flux maps plotted for the two distributor plates at $U_f/U_{mf,2D} = 1.46$ using glass 550 μm particles in Figures 18a, b. The horizontal velocity flux maps show a more uniform mixing for Type A1 distributor plate. For Type B1 distributor plate, the mixing was not

existent in the regions between the adjacent jets. The difference in size of this dead zone region is responsible for higher circulation rates in Type A1 distributor plate when compared with Type A2 and B1 distributor plates.

José et al.²⁴ observed a decrease in circulation rate for increase in orifice diameter, similar to the results observed in the current study. However, Thorley et al.¹ reported the contrary in their experiments. It must be noted that in both studies a single jet was used. Hence, the total orifice flow area as well as the orifice velocity were not sufficient to differentiate the expected solid circulation.

The leveling off of solid circulation rate was observed to occur earlier for Type B1 distributor plate, as shown in Figure 17. At the maximum fluidization velocity ($U_f/U_{mf,2D} = 1.68$), the effect of the jet bypassing the bed height increases the solid circulation rate for Type B1 distributor plate.

Conclusions

PIV was used to investigate the evolution of multiple inlet gas jets located at the distributor base of a two-dimensional fluidized bed setup. A correlation was developed to estimate the solid volume fraction using digital image analysis on the results. The solid volume fraction profile and PIV results were used to quantify the solid circulation rate in the fluidized bed. The effects of fluidization velocity, orifice diameter, orifice pitch, particle diameter, and particle density were studied on the solid circulation rate. The solid volume fraction model compared well with a similar model suggested by Link et al.^{36,37} For all test cases, the mass of bed media estimated using this model was within 15% of the actual bed mass. The solid volume fraction profiles were similar to those observed in the literature.

Solid circulation was found to increase linearly with an increase in fluidization velocity until the jet systems transitioned from isolated jets to interacting systems. For an interacting jet system, the solid circulation increased with fluidization velocity but at a much lower rate. In a multiple jet system, this phenomenon may indicate an optimum operating condition for a bed that requires high circulation rate for low air input. The accuracy of the solid circulation rate calculated in the current experimental study was verified by the circulation rate calculated by the counter-current flow model in the literature.^{8–10}

The circulation rate was observed to slightly increase with the decrease in particle density from 2500 kg/m^3 to 1079 kg/m^3 at same $U_f/U_{mf,2D}$. However, the horizontal particle entrainment velocities decreased significantly with this change. For $U_f/U_{mf,2D} > 1$, the solid circulation rate and entrainment velocities for glass 750 μm particles exceeded that of glass 550 μm particles.

At the same orifice velocities, a higher circulation rate was noted for distributors with smaller orifice pitch. At the same fluidization velocities and for distributors with different orifice diameters but the same orifice flow area, the circulation rate was noted to be much higher for distributors with smaller diameter orifices. Hence, for a distributor plate design, the plate with large number of small diameter orifices was found to enhance the solid circulation rate.

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Notation

- A = constant to convert I_{cor} to area fraction projected at front wall
 A_d = area of distributor, m^2
 B = constant to convert projected area fraction to ε_s
 D_j = jet diameter, m
 d_o = diameter of orifice in distributor plate, m
 d_p = mean particle diameter, m
 f_{neg} = fraction of bed moving with a negative vertical velocity at a specific fluidization velocity
 I_{cor} = corrected image intensity averaged over an interrogation area
 I_m = manipulated intensity of the captured image
 L_j = jet penetration length, m
 \dot{m}_x' = horizontal particle mass flux in the interrogation area; $= \rho_p \varepsilon_s \bar{u}$, $\text{kg/m}^2 \text{ s}$
 \dot{m}_y' = vertical particle mass flux in the interrogation area; $= \rho_p \varepsilon_s \bar{v}$, $\text{kg/m}^2 \text{ s}$
 $\dot{m}_{x,j}$ = horizontal particle entrainment rate of a particular jet, kg/s
 $\dot{m}_{x,\text{tot}}$ = total solid circulation rate of the bed (kg/s), $= \sum_{j=1}^N \dot{m}_{x,j}$
 N = number of orifices within the distributor plate
 P = pitch of orifices within the distributor, m
 U_{avg} = average of \bar{u} calculated at jet periphery for all jets in a bed at a set air flow rate
 U_f = superficial gas velocity, m/s
 U_{mf} = superficial gas velocity at minimum fluidization, m/s
 $U_{\text{mf},2\text{D}}$ = experimental U_{mf} of corresponding distributor plate and particle test case, m/s
 $U_{\text{mf},3\text{D}}$ = experimental minimum fluidization velocity measured in a cylindrical bed, $= U_{\text{mf}}$, m/s
 U_o = mean orifice velocity (m/s), $= (U_f A_d)/N(\pi d_o^2/4)$
 \bar{u} = horizontal particle velocity vector in the interrogation area, m/s
 \bar{v} = vertical particle velocity vector in the interrogation area, m/s
 $\bar{v}_{\text{neg,avg}}$ = mean negative vertical velocity of the particles at specific fluidization velocity
 x = horizontal location along the width of fluidized bed, m
 y = vertical location from the distributor base, m

Greek letters

- α = particle entrainment angle, $^\circ$
 ε_f = mean void fraction in the interrogation area along the bed thickness, $= 1 - \varepsilon_s$
 ε_s = mean solid volume fraction in the interrogation area along the bed thickness
 $\varepsilon_{s,b}$ = bulk solid volume fraction calculated along the expanded bed dimensions
 ρ_p = particle density, kg/m^3

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